



\* From Nowacki (1967).

Although no specific examples of potentially useful triple Bragg reflections are presently known, some progress has been made in narrowing the range of possibilities to be considered. The conditions on the structure amplitudes [equation (2)] are quite restrictive. Let us confine our attention to forbidden reflections that are forbidden strictly, that is, by virtue of space-group symmetry *(International Tables for X-ray Crystallography,* 1969). Then it is not difficult to show that only 5 space groups (of unfortunately infrequent occurrence) among the 10 belonging to class *m3m* of the cubic system can satisfy equation (2). Similarly, only 2 of the 8 most common (for inorganic crystals) space groups (Nowacki, 1967) are eligible. Unfortunately the cubic diamond structure (space group *Fd3m*) and the hexagonal close-packed structure (space group *P63/mmc),* both of which are suitable for double

Bragg reflection (Kottwitz, 1968), are ruled out by the structure-amplitude conditions. For these 7 eligible space groups a limited amount of trial-and-error calculation has been done to find sets of reflections that satisfy equations (1) and (2). For each space group, a sample set of such reflections is given in Table 1, together with the crystal system and the number of known crystals (Nowacki, 1967).

### **References**

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# Polarization factor for graphite X-ray monochromators. By HÅKON HOPE, *Department of Chemistry, University of California, Davis, California* 95616, *U.S.A.*

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Comparison of Ni-filtered and graphite-monochromatized Cu  $K\alpha$  diffraction data indicates that the 'ideally imperfect' monochromator polarization factor is not applicable.

Since graphite was introduced as an X-ray monochromator the assumption of the polarization factor being that of an ideally imperfect crystal appears to have been tacitly accepted. A communication by Miyake & Togawa (1964), however, clearly points out that for high precision work it becomes necessary to test the validity of this assumption.

After a graphite monochromator had been installed on the Picker diffractometer in this laboratory part of the initial testing consisted of a comparison of a Cu *Ka* (Ni filter) and a Cu  $K\alpha$  (graphite monochromator) data set from the same crystal. For the comparison to be valid the specimen crystal must be perfectly stable towards X-ray irradiation, and secondary extinction effects must be small. A filter data set  $(F_t)$  satisfying these criteria was available from a reinvestigation of  $p, p'$ -dichloroazobenzene (initial study by Hope & Victor, 1969.) The structure had been refined to  $R=0.032$ , with  $|F_{\rho}|$  being about 5% smaller than  $|F_c|$  for the most intense reflection. The crystal had been removed from the diffractometer with the goniometer head attached to the mounting base, and when the assembly was remounted on the goniostat it was found that the original setting parameters could be used without change.

The Cu  $K\alpha$  (monochromator) data set was collected with exactly the same settings, scan ranges and backgroundcount times as the filter data. About 950 reflections with  $I_{\alpha}$  > 2 $\sigma(I_{\alpha})$  which were common for the two data sets were used for comparison of polarization factors. The monochromator intensities were first reduced to F's *(FM)* by use of the polarization factor for an ideally imperfect monochromator with diffraction vector in a plane normal to the diffractometer equator, given by  $p_K = (\cos^2 2\theta + \cos^2 2\theta_M)/$  $(1 + \cos^2 2\theta_M)$ , where  $\theta$  is the Bragg angle for the reflection under consideration and  $\theta_M$  the monochromator Bragg angle (Azároff, 1955). [The corresponding expression for the 'perfect' monochromator is  $p_D = (\cos^2 2\theta + |\cos 2\theta_M|)/$  $(1 + |\cos 2\theta_M|).$ 

Average normalized  $F_M/F_f$  ratios were calculated for 5<sup> $\circ$ </sup> ranges in  $\theta$  with the results plotted in Fig. 1. Although small, the hump centered at  $\theta = 45^\circ$  clearly points to a systematic, angle dependent error.

Following the procedure suggested by Miyake & Togawa (1964) the monochromator polarization factor was then expressed as  $p = cpD+(1-c)pK$ . From the shape of the initial  $F_M/F_f$  curve a value of  $c = 0.65$  was estimated, and a new set of  $F_M$  was calculated. The average ratio was plotted as before, with the result also shown in Fig. 1. We see that above  $\theta \sim 35^{\circ}$  the observed ratios very closely approach unity. At lower  $\theta$  values there is a gradual drop to a ratio of 0.99 at  $15^{\circ}$ , pointing to some other systematic error. Inspection of the raw filter data revealed some degree of



Fig. 1. Ratios  $F_M/F_f$  against  $\theta$ . Dots represent ratios derived by use of  $p<sub>K</sub>$ , and triangles show ratios resulting from  $p=$  $0.65p_D + 0.35p_K$ . The curve gives  $(p_K/(0.65p_D + 0.35p_K))^{1/2}$  as a function of  $\theta$ .

skewness in background at lower  $\theta$  angles, in part resulting from background measurements being made at points where 'white' radiation is strongly absorbed by the filter, giving rise to underestimated background values.

The possibility of an intensity-related error was ruled out by comparing  $F_M$  with  $F_f$  as a function of F, with no systematic trend apparent. An R index of  $0.010$  ( $R=2$ -  $(\Sigma |F_M-F_f|)/\Sigma(F_M+F_f)$  calculated for the reflections used indicates a very satisfactory overall agreement between the two data sets.

The results obtained in this study show that commercially available graphite monochromators can behave quite differently from 'ideally imperfect' crystals, and that allowance should be made for any departure from ideal behavior. Each monochromator must of course be calibrated; it is also conceivable that the calibration might change as a result of irradiation.

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### **Reference**

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